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A COMPUTER SUBROUTINE FOR STRESS ANALYSIS OF
ROTATING, HEATED DISKS

by

John E. Brock

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TABLE OF CONTENTS

| | |
|---|----|
| Introduction----- | 1 |
| Fundamental Analysis----- | 2 |
| Power Law of Thickness Variation----- | 3 |
| Exponential Law of Thickness Variation----- | 4 |
| General Law of Thickness Variation----- | 5 |
| Computer Implementation----- | 8 |
| References----- | 9 |
| Appendix A----- | 10 |
| Appendix B----- | 13 |
| Appendix C----- | 16 |
| Initial Distribution List----- | 20 |

A Computer Subroutine for Stress
Analysis of Rotating, Heated Disks
by
John E. Brock and Robert E. Brown

Introduction

Although, as is indicated by the title hereof, the principal purpose of this monograph is to present tested and proved digital computer software for the analysis of stress in a spinning axisymmetric disk having a radially variable thermal strain field, the opportunity is also taken of developing the theory and presenting some analytic solutions.

The method developed herein for computer analysis of disks having a general law of thickness variation was suggested by the algorithm contained in reference 2 and it appears to have advantages over such procedures as that of M. Donath, reference 3, which has been widely circulated in a book by S. Timoshenko, reference 5.

In what follows we immediately obtain a second order linear differential equation with dependent variable u , the radial deformation, and r , the radius. Analytic treatment is given for two particular laws of thickness variation. For the general case of thickness variation, the equation is recast as a second order linear differential equation in which the dependent variable is radial stress, σ_r . However, for numerical treatment an alternate form is more useful and direct and this forms the basis of the digital computer software which is given

and illustrated in the appendices hereto.

Fundamental Analysis

We presume that the disk is thin enough and that the thickness varies slowly enough with respect to radius that we are justified in neglecting all stress components excepting only the radial stress σ_r and the circumferential stress σ_θ . Material properties E , Young's modulus, and ν , Poisson's ratio, are presumed to be indeed constant. The thermal strain field, αT , and the density γ may be specified functions of radius.

Two types of problem are considered:

1. Annular disk, $0 < a \leq r \leq b$, with $\sigma_r(a)$ and $\sigma_r(b)$ being specified.
2. Solid disk, $0 \leq r \leq b$, with $\sigma_r(b)$ being specified.

We also use the symbols $t = t(r)$ for disk thickness and ω for angular velocity. Other simplifying notation will be introduced later on.

Consideration of radial equilibrium leads without difficulty to the equation

$$\frac{1}{t} \frac{d(rt\sigma_r)}{dr} = \sigma_\theta - \gamma\omega^2 r^2 \quad (1)$$

The thermoelastic constitutive equations are

$$E\varepsilon_\theta = \sigma_\theta - \nu\sigma_r + E\alpha T; \quad E\varepsilon_r = \sigma_r - \nu\sigma_\theta + E\alpha T \quad (2a,b)$$

where the strain components are

$$\varepsilon_\theta = u/r; \quad \varepsilon_r = du/dr \quad (3a,b)$$

Strain compatibility leads to the equation

$$r \frac{d}{dr}(Eu/r) = -(1+\nu)(\sigma_\theta - \sigma_r) \quad (4)$$

These equations may easily be combined into the differential equation

$$\begin{aligned}
u'' + u'/r - u/r^2 + (t'/t)[u' + vu/r - (1+v)\alpha T] \\
= (1+v)(\alpha T)' - (1-v^2)\gamma\omega^2 r/E
\end{aligned} \tag{5}$$

where primes denote differentiation with respect to r .

Two particular laws of thickness variation permit simple analytic treatment.

Power Law of Thickness Variation

If

$$t = t_0 (r/r_0)^n \tag{6}$$

so that

$$(t'/t) = n/r \tag{7}$$

then equation 5 becomes

$$u'' + (1+vn)u'/r + (vn-1)u/r^2 = \beta' + n\beta/r - kr \tag{8}$$

where

$$\beta = (1+v)\alpha T, \quad k = (1-v^2)\gamma\omega^2/E \tag{9,10}$$

Equation 8 may be rewritten as

$$\frac{d}{dr} \left[r^{1-m} \frac{d}{dr} (r^{(m+n)/2} u) \right] = r^{(2+n-m)/2} (\beta' + n\beta/r - kr) \tag{11}$$

where

$$m = \pm \sqrt{(n^2 - 4n + 4)} = \pm \sqrt{(n-2)^2 + 4(1-n)} \tag{12}$$

and either the positive or the negative sign may be used. Equation 11 may be proved simply by performing the indicated operations and comparing with equation 8.

The quantity on the right in equation 11 is well defined so that the solution of the differential equation may be obtained simply by integration, multiplication by r^{m-1} , and another integration. Two constants of integration are introduced. For the solid disk (case 2)

$u(0) = 0$ gives one condition and the second comes from satisfying the given value of $\sigma_r(b)$. For the annular disk (case 1) satisfying the given conditions $\sigma_r(a)$ and $\sigma_r(b)$ permits evaluating the constants.

An example with $n = -0.42$ is given in Appendix C. Note that the case $n = 0$ corresponds to a disk of uniform thickness.

Exponential Law of Thickness Variation

If

$$t = t_0 \exp(-mr^2) \quad (13)$$

where m is a constant of appropriate dimensionality, then

$$(t'/t) = -2rm \quad (14)$$

and equation 5 becomes

$$u'' + (1/r - 2rm)u' - (1/r^2 + 2vm) = -2rm\beta + \beta' - kr \quad (15)$$

If additionally we assume that $\beta' = 0$ (which makes the thermal strain field constant — a triviality) the solution is simply

$$u = r(\beta + k/2m)/(1+v) \quad (16)$$

In this case we can easily find

$$\sigma_r = \sigma_\theta = \gamma\omega^2/2m \quad (17)$$

which is independent of r . Thus, if an allowable normal stress σ_A is specified and if blade or bucket loading at $r = b$ is w (pounds, say) per unit circumference, then a disk having thickness

$$t = (w/\sigma_A) \exp[(b^2 - r^2)(\gamma\omega^2/2\sigma_A)] \quad (18)$$

will be such that $\sigma_r \equiv \sigma_\theta \equiv \sigma_A$. If the failure criterion is the maximum shearing stress criterion (Tresca's condition), it is clear that this disk is optimal in the sense of having least volume and

thus having least weight. If the failure criterion were that of von Mises, a slightly lighter disk would suffice.

General Law of Thickness Variation

Equations 3a and 2a give

$$Eu = rE\alpha T + r(\sigma_{\theta} - v\sigma_r) \quad (19)$$

and by differentiation

$$Eu' = E\alpha T + rE(\alpha T)' + (\sigma_{\theta} - v\sigma_r)' + r(\sigma_{\theta}' - v\sigma_r') \quad (20)$$

From 3b and 2b we also have

$$Eu' = E\alpha T + \sigma_r - v\sigma_{\theta} \quad (21)$$

Subtracting 21 from 20 and rearranging gives

$$E(\alpha T)' + (1+v)(\sigma_{\theta} - \sigma_r)/r + \sigma_{\theta}' - v\sigma_r' = 0 \quad (22)$$

Equation 1 may be rewritten as

$$\sigma_{\theta} - \sigma_r = r\sigma_r' + rv\sigma_r' + \gamma\omega^2 r^2 \quad (23)$$

where we have written

$$v = t'/t \quad (24)$$

for convenience. Differentiating 23 we get

$$\sigma_{\theta}' = 2\sigma_r' + r\sigma_r'' + v\sigma_r + rv\sigma_r' + rv'\sigma_r + 2\gamma\omega^2 r \quad (25)$$

and substituting 23 and 25 into 22 gives

$$r\sigma_r'' + (3+rv)\sigma_r' + [(2+v)v+rv']\sigma_r + (3+v)\gamma\omega^2 r + E(\alpha T)' = 0 \quad (26)$$

This is a single differential equation with dependent variable σ_r and can be dealt with by standard numerical methods. The conditions for evaluating the constants of integration have been mentioned earlier.

However, the preceding procedure is not particularly satisfactory. For one thing, the solid disk case for which r can become

zero encounters numerical difficulties unless special treatment is employed to evade them. More importantly, however, if T is given by a graph or a numerical table, determination of $(\alpha T)'$ may involve numerical difficulties which, ultimately, are due to our having performed the differentiation to arrive at equation 20. Accordingly, an alternate procedure which adheres more closely to the fundamental mechanics of the problem is described in what follows.

We consider the annular disk first and we represent the unknown stress difference $\sigma_\theta - \sigma_r$ in the form

$$\sigma_\theta - \sigma_r = (A + B\eta)r \quad (27)$$

where A and B are unknown constants and η is an unknown function normalized so that

$$\eta(a) = 0, \quad \eta(b) = 1 \quad (28a,b)$$

Initially we make an assumption for η taking a linear variation in the absence of better information. By the use of some of the preceding equations we will be able to construct an improved form for η and will iterate until there is satisfactory convergence.

Let

$$z = Eu/r, \quad w = ExT, \quad \beta = \gamma\omega^2 \quad (29,30,31)$$

noting that w and β have different meanings here than when they were used earlier. We can recast equation 1 as

$$d(t\sigma_r)/dr = (\sigma_\theta - \sigma_r)t/r - \gamma\omega^2 rt = t(A + B\eta - \gamma\omega^2 r) \quad (32)$$

Before performing the indicated integration we introduce two convenient notational devices, viz.

$$p... = \int_0^r ... dr; \quad *... = [...]_{r=b} \quad (33,34)$$

Thus, from equation 32 we have

$$t\sigma_r = (t\sigma_r)_a + A\rho t + B\rho nt - p\beta rt \quad (35)$$

From equation 4 and equation 27 we have

$$dz/dr = -(1+\nu)(A+B\eta) \quad (36)$$

so that

$$z = (z)_a - (1+\nu)[A(r-a) + B\eta] \quad (37)$$

Equation 2a is

$$z = w + (\sigma_\theta - \sigma_r) + (1-\nu)\sigma_r = w + (A+B\eta)r + (1-\nu)\sigma_r \quad (38)$$

so that

$$(z)_a = (w)_a + Aa + (1-\nu)(\sigma_r)_a \quad (39)$$

$$(z)_b = (w)_b + (A+B)b + (1-\nu)(\sigma_r)_b \quad (40)$$

Evaluating 38 at $r = b$ and using 39 and 40 gives one equation involving the unknowns A and B. Evaluating 35 at $r = b$ gives a second. These equations can be arranged as

$$\begin{bmatrix} \rho(pt) & \rho(pnt) \\ (2+\nu)(b-a) & b+(1+\nu)\rho(p\eta) \end{bmatrix} \begin{bmatrix} A \\ B \end{bmatrix} = \begin{bmatrix} \rho(p\beta rt) + (t\sigma_r)_b - (t\sigma_r)_a \\ (w)_a - (w)_b + (1-\nu)[(\sigma_r)_a - (\sigma_r)_b] \end{bmatrix} \quad (41)$$

so that one can easily solve for A and B. Then σ_r is obtained from 35, z is obtained from 37 and 39, and $(\sigma_\theta - \sigma_r)$ and σ_θ are obtained from 38. Then a new function η is calculated from

$$\eta = [(\sigma_\theta - \sigma_r)/r - A]/B \quad (42)$$

Using the new η the process is iterated, convergence being monitored by examination of the sequence of values of A and B that are calculated. When convergence is satisfactory, the desired functions σ_r and σ_θ are at hand.

The situation is simpler for the solid disk, case 2. $(\sigma_r)_a$ is not given but conditions of continuity require that $(\sigma_\theta - \sigma_r)/r$ vanish at $r = 0$. Thus A is zero and B can be obtained from

$$B = \frac{(w)_o - (w)_b + c[1-(t)_o/(t)_b](t\sigma_r)_b + c^*(p\beta r t)}{b + (1+v)^*(p\eta) + c^*(p\eta t)} \quad (43)$$

where

$$c = (1-v)/(t)_o \quad (44)$$

Otherwise the procedure is as for case 1.

Computer Implementation

The theory embodied in equations 27 through 44 and the associated procedure has been programmed for digital computer using the FORTRAN language. An initial programming based directly on the preceding equations and written in January 1978 by the junior author hereof has been supplanted by a newer programming which is somewhat more compact due to the employment therein of ancillary subroutines developed for use in another problem (the lateral buckling of elastic beams) on which we are working. This program, actually a subroutine called RODISK is listed in Appendix A hereof. This listing itself contains comments which adequately explain the construction of a MAIN program which supplies necessary input information and which invokes RODISK. Appendix B lists the ancillary subroutines. In each case comments indicate the purpose and employment of the subroutine. These may prove useful in constructing the MAIN program. For this reason, the ancillary subroutines DUPV and PRIV are given even though they are not called by RODISK.

In the computer implementation, the various *functions* of r which appear in the theory are represented by *vectors* the elements of which are function values at equally spaced values of r .

The ancillary subroutines manipulate these vectors. All of these are obvious except possibly INTV which performs an integration by use of Milne's formulas, cf. reference 4.

In the subroutine RODISK there is a slight departure from the theory as given herein. As a first step, all quantities and functions were "dedimensionalized" but otherwise the method is just as described above. Somewhat finer details of what was actually done may be found in reference 1.

Appendix C contains some examples and remarks concerning them.

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3. Donath, M., *Die Berechnung rotierender Scheiben und Ringe*, Berlin, 1912.
4. Milne, W. E., *Numerical Calulus*, Princeton Univ. Press, 1949, p. 119.
5. Timoshenko, S. *Strength of Materials, Part II*, 3rd ed., D. van Nostrand Co., Inc., 1956, pp. 223 - 228.

Appendix A

Listing of subroutine RODISK

(The listing on this page, page 10, is of the comments which provide instructions for the use of RODISK. Commands appear on the following two pages.)

SUBROUTINE RODISK. JOHN E. BROCK., 1 MAY 1978
THIS IS A SUBROUTINE FOR DETERMINING RADIAL AND CIRCUMFERENTIAL STRESSES IN AN AXISYMMETRIC THIN ELASTIC DISK, ROTATING AT ANGULAR VELOCITY CMEGA (RADTANS PER SECCND) ABOUT ITS AXIS OF SYMMETRY AND HAVING AN AXISYMMETRIC DISTRIBUTION OF THERMAL STRAIN. TWO TYPES OF PROBLEM MAY BE TREATED:
TYPE 1. ANNULAR DISK OF INSIDE RADIUS ARAD AND OUTSIDE RADIUS BRAD. THE RADIAL STRESS IS SRA AT THE INNER RADIUS AND SRB AT THE OUTER RADIUS. THE INSIDE RADIUS MUST BE GREATER THAN ZERO.
TYPE 2. SOLID DISK OF OUTSIDE RADIUS BRAD AND WITH RADIAL STRESS SRB AT THE OUTSIDE RADIUS.
THE USER MUST PROVIDE A MAIN PROGRAM WHICH CALLS SUBROUTINE RODISK AFTER IT HAS SUPPLIED THE FOLLOWING INFORMATION.
(1) N, INTEGER. (N-1) IS THE NUMBER OF EQUAL SUBDIVISIONS INTO WHICH THE ANNULAR RADIUS (BRAD MINUS ARAD) IS DIVIDED FOR COMPUTATIONAL PURPOSES. THE PRESENT DIMENSIONING CAN ACCOMMODATE N NOT LARGER THAN 101.
(2) BRAD
(3) ARAD (NOT NECESSARY FOR PROBLEMS OF TYPE 2.)
(4) SRB
(5) SRA (NOT NECESSARY FOR PROBLEMS OF TYPE 2.)
(6) TEEBEE, DISK THICKNESS AT OUTSIDE RADIUS
(7) POIS, POISSON'S RATIO.
(8) KP(1) = 1,2. INTEGER TO DENOTE PROBLEM OF TYPE 1,2.
(9) KP(2), INTEGER TO PROVIDE FOR SKIPPING WHILE PRINTING OUTPUT. FOR EXAMPLE, IF N=101 AND KP(2)=5, ONLY EVERY FIFTH SET OF VALUES WILL BE PRINTED: 1ST, 6TH, 11TH, 16TH, 21ST, 26TH, 31ST, 36TH, 41ST, 46TH, 51ST, 56TH, 61ST, 66TH, 71ST, 76TH, 81ST, 86TH, 91ST, 96TH, 101ST.
(10) KP(3), INTEGER SPECIFYING THE NUMBER OF ITERATIONS TO BE PERFORMED. USUALLY KP(3)=10 IS SUFFICIENT FOR ENGINEERING ACCURACY.
(11) KP(4). IF KP(4)=0 ONLY FINAL ANSWERS WILL BE PRINTED IF KP(4)=1 A SEQUENCE OF ITERANT VALUES WILL BE PRINTED, INDICATING DEGREE OF CONVERGENCE.
(12) VECTORS X(I,J), I=1,2,3; J=1,2,...,N.
VECTOR X(1,J) CONTAINS VALUES OF THE RATIO (LOCAL THICKNESS OF DISK)/(TEEBEE) COMPUTED AT EQUALLY SPACED RADII STARTING AT THE INSIDE AND ENDING AT THE OUTSIDE.
VECTOR X(2,J) CONTAINS VALUES OF (GAMMA)(OMEGA-SQUARED) TIMES (BRAD-SQUARED) DIVIDED BY (SRB). FOR MOST PROBLEMS GAMMA DOES NOT VARY WITH RADIUS AND THIS QUANTITY IS A CONSTANT.
VECTOR X(3,J) CONTAINS VALUES OF (E)(ALPHA)(TEE)/(SRB) WHERE (E) IS YOUNG'S MODULUS, (ALPHA) IS THE COEFFICIENT OF LINEAR THERMAL EXPANSION, AND (TEE) IS TEMPERATURE CHANGE.
THE MAIN PROGRAM MUST CONTAIN THE STATEMENTS:
IMPLICIT REAL*8 (A-H,O-Z)
INTEGER KP(4)
COMMON X,N
COMMON /ONE/ARAD,BRAD,SRA,SRB,TEEBEE,POIS,KP
FOLLOWING SUBROUTINE RODISK THERE ARE SEVERAL ANCILLARY SUBROUTINES WHICH PERFORM VARIOUS OPERATIONS ON THE VECTORS X(I,J). THE FUNCTION OF EACH IS OBVIOUS FROM THE LISTING. THEY MAY BE USED IN THE USER'S MAIN PROGRAM. TWO OF THESE ANCILLARY SUBROUTINES WHICH ARE AVAILABLE IN THIS PACKAGE BUT WHICH ARE NOT CALLED BY SUBROUTINE RODISK ARE DUPV WHICH DUPLICATES A VECTOR AND PRIV WHICH PRINTS A VECTOR.

```

SUBROUTINE RODISK
IMPLICIT REAL*8(A-H,O-Z)
REAL*8 X(20,101)
INTEGER KP(4)
COMMON /ONE/ARAD,BRAD,SRA,SRB,TEEBEE,POIS,KP
ONE=1.D+0
POIS=3.D-1
ZERO=C.D+0
IF(KP(1).EQ.2) ARAD=ZERO
RHC=ARAD/BRAD
ENM=N-1
WRITE(6,2) KP(1)
2 FORMAT(/,10X,'RODISK PROBLEM OF TYPE ',I1,'.',//)
DO 5 I=1,N
EIM=I-1
Y=EIM/ENM
X(4,I)=RHO+(ONE-RHO)*Y
X(5,I)=Y
5 X(6,I)=Y
ITER=1
ETA=X(1,1)/X(1,N)
IF(KP(1).EQ.2) GO TO 100
C THE PROBLEM IS OF TYPE ONE: ANNULAR DISK
SRAT=SRA/SRB
C1=(2.D+0+POIS)*(ONE-RHO)
CALL INTV(1,7)
C2=X(7,N)
C5=X(3,1)-X(3,N)-(ONE-POIS)*(ONE-SRAT)
CALL MULV(1,2,8)
CALL MULV(8,4,9)
CALL INTV(9,10)
C6=X(10,N)+(ONE-ETA*SRAT)/(ONE-RHO)
20 CALL INTV(6,11)
C3=ONE+(ONE-RHO)*(ONE+POIS)*X(11,N)
CALL MULV(1,6,12)
CALL INTV(12,13)
C4=X(13,N)
D=C1*C4-C2*C3
A=(C5*C4-C6*C3)/D
B=(C1*C6-C2*C5)/D
IF(KP(4).EQ.1) WRITE(6,7) ITER,A,B
7 FORMAT(5X,I10,1P2E20.5)
CALL MULS(7,14,A)
CALL MULS(13,15,B)
CALL ADDV(14,15,15)
CALL SUBV(15,10,16)
S=CNE-RHO
CALL MULS(16,16,S)
S=ETA*SRAT
CALL ADDS(16,16,S)
CALL DIVV(16,1,20)
ZJ=A*RHO+X(3,1)+(ONE-POIS)*SRAT
CALL MULS(11,11,3)
CALL MULS(5,16,A)
CALL ADDV(11,16,16)
S=-(ONE-RHO)*(ONE+POIS)
CALL MULS(16,16,S)
CALL ADDS(16,19,Z0)
S=POIS-ONE
CALL MULS(20,18,S)
CALL ADDV(18,19,18)
CALL SUBV(18,3,19)
ITER=ITER+1
IF(ITER.GT.KP(3)) GO TO 200
CALL DIVV(19,4,17)
CALL SUBS(17,17,A)
CALL DIVS(17,6,B)
GO TO 20

```

```

100 C1=PCIS-ONE
    C2=ETA
    C5=X(3,1)-X(3,N)+POIS-ONE
    CALL MULV(1,2,7)
    CALL MULV(7,4,8)
    CALL INTV(8,9)
    C6=ONE+X(9,N)
105 CALL INTV(6,10)
    C3=ONE+(ONE+POIS)*X(10,N)
    CALL MULV(1,6,11)
    CALL INTV(11,12)
    C4=X(12,N)
    D=C1*C4-C2*C3
    SRAT=(C5*C4-C6*C3)/D
    B=(C1*C6-C2*C5)/D
    IF(KP(4).EQ.1) WRITE(6,7) ITER,SRAT,B
    CALL MULS(12,13,B)
    CALL SUBV(13,9,14)
    S=ETA*SRAT
    CALL ADDS(14,20,S)
    CALL DIVV(20,1,20)
    S=-B*(ONE+POIS)
    CALL MULS(10,15,S)
    S=POIS-ONE
    CALL MULS(20,18,S)
    S=X(3,1)+(ONE-POIS)*SRAT
    CALL ADDS(18,19,S)
    CALL SUBV(19,3,19)
    S=-B*(ONE+POIS)
    CALL MULS(10,17,S)
    CALL ADDV(19,17,19)
    ITER=ITER+1
    IF(ITER.GT.KP(3)) GO TO 200
    CALL MULS(4,16,B)
    X(16,1)=ONE
    CALL DIVV(19,16,6)
    GO TO 105
200 CALL MULS(4,7,BRAD)
    CALL MULS(20,8,SRB)
    CALL MULS(19,9,SRB)
    CALL ADDV(8,9,9)
    S=SRB/BRAD**2
    CALL MULS(2,10,S)
    CALL MULS(3,11,SRB)
    CALL MULS(1,15,TEEBEE)
    WRITE(6,204)
204 FORMAT(////)
    WRITE(6,205)
205 FORMAT(24X,'RADIUS',11X,'THICKNESS',6X,'GAMMA.OMEGA.SQ'
1'EE.ALPHA.TEE',8X,'SIGMA.RADIAL',7X,'SIGMA.CIRCUMF')
    KSKIP=KP(2)
    DO 210 I=1,N,KSKIP
    J=I/KSKIP
    WRITE(6,211)J,X(7,I),X(15,I),X(1,I),X(11,I),X(8,I),X(9
210 CONTINUE
211 FORMAT(I10,1P6E20.5)
    RETURN
    END

```

Appendix B

Listing of ancillary subroutines

```
C      SUBROUTINE ADDV(N1,N2,N3)
C      ANCILLARY SUBROUTINE
C      ADC TWO VECTORS TERM BY TERM
      X(N3,I)=X(N1,I)+X(N2,I)
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1  X(N3,I)=X(N1,I)+X(N2,I)
      RETURN
      END

C      SUBROUTINE SUBV(N1,N2,N3)
C      ANCILLARY SUBROUTINE
C      SUBTRACT TWO VECTORS TERM BY TERM
      X(N3,I)=X(N1,I)-X(N2,I)
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1  X(N3,I)=X(N1,I)-X(N2,I)
      RETURN
      END

C      SUBROUTINE MULV(N1,N2,N3)
C      ANCILLARY SUBROUTINE
C      MULTIPLY TWO VECTORS TERM BY TERM
      X(N3,I)=X(N1,I)*X(N2,I)
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1  X(N3,I)=X(N1,I)*X(N2,I)
      RETURN
      END

C      SUBROUTINE DIVV(N1,N2,N3)
C      ANCILLARY SUBROUTINE
C      DIVIDE TWO VECTORS TERM BY TERM
      X(N3,I)=X(N1,I)/X(N2,I)
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1  X(N3,I)=X(N1,I)/X(N2,I)
      RETURN
      END

C      SUBROUTINE ADDS(N1,N2,S)
C      ANCILLARY SUBROUTINE
C      ADC A SCALAR TO EACH TERM OF A VECTOR
      X(N2,I)=X(N1,I)+S
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1  X(N2,I)=X(N1,I)+S
      RETURN
      END
```



```

C      SUBROUTINE SUBS(N1,N2,S)
C      ANCILLARY SUBROUTINE
C      SUBTRACT A SCALAR FROM EACH TERM OF A VECTOR
      X(N2,I)=X(N1,I)-S
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1 X(N2,I)=X(N1,I)-S
      RETURN
      END

C      SUBROUTINE MULS(N1,N2,S)
C      ANCILLARY SUBROUTINE
C      MULTIPLY EACH TERM OF A VECTOR BY A SCALAR
      X(N2,I)=X(N1,I)*S
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1 X(N2,I)=X(N1,I)*S
      RETURN
      END

C      SUBROUTINE DIVS(N1,N2,S)
C      ANCILLARY SUBROUTINE
C      DIVIDE EACH TERM OF A VECTOR BY A SCALAR
      X(N2,I)=X(N1,I)/S
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1 X(N2,I)=X(N1,I)/S
      RETURN
      END

C      SUBROUTINE DUPV(N1,N2)
C      ANCILLARY SUBROUTINE
C      DUPLICATE A VECTOR
      X(N2,I)=X(N1,I)
      REAL*8 X(20,101),S
      COMMON X,N
      DO 1 I=1,N
1 X(N2,I)=X(N1,I)
      RETURN
      END

C      SUBROUTINE INTV(N1,N2)
C      ANCILLARY SUBROUTINE
C      INTEGRATE A VECTOR USING MILNE'S METHOD
      REAL*8 X(20,101),EN,R,ADD,NINO,NTNO,FIVO,THTO
      COMMON X,N
      EN=N-1
      EN=1.D+0/EN
      R=EN/2.4D+1
      NINO=R*9.D+0
      NTNO=R*1.9D+1
      FIVO=R*5.D+0
      THTO=R*1.3D+1
      X(N2,1)=0.D+0
      X(N2,2)=NINO*X(N1,1)+NTNO*X(N1,2)-FIVO*X(N1,3)+R*X(N1,4)
      NM3=N-3
      DO 1 K=1,NM3
      KP1=K+1
      KP2=K+2
      KP3=K+3
      ADD=THTO*(X(N1,KP1)+X(N1,KP2))-R*(X(N1,K)+X(N1,KP3))
1 X(N2,KP2)=X(N2,KP1)+ADD
      X(N2,N)=X(N2,N-1)+NINO*X(N1,N)+NTNO*X(N1,N-1)-FIVO*X(N1,N-1)+R*X(N1,N-3)
      RETURN
      END

```


Appendix C

Examples

The listing on the next page is of a MAIN program which calls RODISK twice to solve two different problems. On an IBM 360/67 compile time (for the MAIN, RODISK, and the ancillary subroutines) was 12 s, link time was 2 s, and execution time for both problems was 1.5 s. In both problems we used $N = 41$ and iterated 10 times.

The first problem is of type 2 (solid disk) with $b = 10$, $t = 1/(1.6 + .008r^2)$, $\nu = 0.3$, $\gamma\omega^2 = 120$, $\sigma_r(b) = 14000$, and $E\alpha T = 25105 + 1300 \log_e(t) - r^2(233 + 16t)$ Units are inches and pounds. The problem was made up from the exact solution

$$\sigma_r = 9000 + 50r^2; \quad \sigma_\theta = \sigma_r + r^2(120 + 16t)$$

The results, shown on page 18, show evaluations for the stress components which are correct within 0.03 psi even though convergence was complete to only about 4 digits as indicated by the sequence of values above the final tabulation; these are values of $\sigma_r(0)/\sigma_r(b)$ and of B .

The second problem is of type 1 with $a = 2.57$, $b = 5.15$, $\sigma_r(a) = 18205$, $\sigma_r(b) = 22000$, $t = 0.1493 r^{-.42}$, $T = 60 - 1.6r^2$, $\nu = 0.3$, and $E\alpha = 194.3$. The units are inches and pounds. Since the thickness variation is a power law, the theoretical solution given in the body hereof may be used to obtain the theoretically exact solution

$$\begin{aligned}\sigma_r &= -113.95r^2 + 15832r^D - 11708r^Q \\ \sigma_\theta &= 122.80r^2 + 13801r^D + 14310r^Q\end{aligned}$$

```

C      J. E. BROCK      1 MAY 1976
C      THIS IS A MAIN PROGRAM TO TEST MY SUBROUTINE RODISK
      IMPLICIT REAL*8(A-H,O-Z)
      REAL*8 X(20,101)
      INTEGER KP(4)
      COMMON X,N
      COMMON /ONE/ARAD,BRAD,SRA,SRB,TEEBEE,POIS,KP
C      THE FIRST PROBLEM IS FOR A SOLID DISK
      KP(1)=2
      KP(2)=4
      KP(3)=10
      KP(4)=1
      BRAD=1.0+1
      SRB=1.40+4
      N=41
      ENM=N-1
      DO 20 I=1,N
      EIM=I-1
      R=EIM*BRAD/ENM
      THICK=1.0+0/(1.60+0+8.0-3*R**2)
      X(2,I)=THICK
      W=2.51050+4+1.30+3*DLOG(THICK)-(R**2)*(2.330+2+1.60+1*THICK
20  X(3,I)=W/SRB
      CONTINUE
      S=X(20,N)
      TEEBEE=S
      CALL DIVS(20,1,S)
      CALL RODISK
C      THE SECOND PROBLEM IS FOR AN ANNULAR DISK
      TEEBEE=1.4930-1*5.150+0*(-4.20-1)
      ENM=N-1
      KP(1)=1
      ARAD=2.570+0
      BRAD=5.150+0
      SRA=1.82050+4
      SRB=2.20+4
      BETA=5.024734710-1
      DO 30 I=1,N
      EIM=I-1
      Y=EIM/ENM
      X(5,I)=Y
      X(6,I)=Y
      R=ARAD+(BRAD-ARAD)*Y
      X(4,I)=R/BRAD
      THICK=1.4930-1*R*(-4.20-1)
      X(1,I)=THICK/TEEBEE
      X(2,I)=BETA
      X(3,I)=(6.0+1-1.60+0*R**2)*1.9430+2/SRB
30  CONTINUE
      CALL RODISK
      STOP
      END

```


where $p = .29171$ and $q = -1.87171$. The computer results agree with five digit accuracy even though the sequence of iterants (A,B) shows only four digit convergence.

The same two problems were also worked with $N = 101$ and 14 iterations. The execution time went from 1.5 s to 3.0 s. The maximum change in any stress value was 0.3 psi. From these and other problems it may be concluded that there are no difficulties of accuracy, computer storage, or execution time.

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